## Strong Instability of Standing Waves for a Type of Hartree Equations

WANG Chenglin<sup>1</sup> and ZHANG Jian<sup>2,\*</sup>

Received 9 September 2022; Accepted 24 September 2024

**Abstract.** In this paper, we study the following three-dimensional Schrödinger equation with combined Hartree-type and power-type nonlinearities

$$i\partial_t \psi + \Delta \psi + (|x|^{-2} * |\psi|^2) \psi + |\psi|^{p-1} \psi = 0$$

with  $1 . Using standard variational arguments, the existence of ground state solutions is obtained. And then we prove that when <math>p \ge 3$ , the standing wave solution  $e^{i\omega t}u_{\omega}(x)$  is strongly unstable for the frequency  $\omega > 0$ .

AMS Subject Classifications: 35Q55, 35B35, 35A15

Chinese Library Classifications: O175.27

Key Words: Hartree equation; standing wave; variational arguments; strong instability; blowup.

## 1 Introduction

Consider the following three-dimensional Schrödinger equation with combined Hartree-type and power-type nonlinearities

$$i\partial_t \psi + \Delta \psi + (|x|^{-2} * |\psi|^2) \psi + |\psi|^{p-1} \psi = 0$$
(1.1)

with  $1 , where <math>\psi = \psi(t, x)$  is a complex-valued wave function in time-space  $(t, x) \in \mathbb{R} \times \mathbb{R}^3$ ; i is the imaginary unit;  $\Delta$  is the Laplace operator on  $\mathbb{R}^3$ ; \* denotes the standard convolution of the integral kernal  $|x|^{-2}$  and the square term  $|\psi|^2$  on  $\mathbb{R}^3$ . More precisely,

$$(|x|^{-2} * |\psi|^2)(x) = \int_{\mathbb{R}^3} \frac{|\psi(y)|^2}{|x - y|^2} dy.$$

<sup>&</sup>lt;sup>1</sup> School of Sciences, Xihua University, Chengdu 610039, China;

<sup>&</sup>lt;sup>2</sup> School of Mathematical Sciences, University of Electronic Science and Technology of China, Chengdu 611731, China.

<sup>\*</sup>Corresponding author. Email addresses: zhangjian@uestc.edu.cn (J. Zhang)

The term  $(|x|^{-2}*|\psi|^2)\psi$  is called the Hartree-type (or the nonlocal nonlinearity) and  $|\psi|^{p-1}\psi$  is called power-type (or the local nonlinearity).

(1.1) origins from the Schrödinger-Poisson-Slater system, which was introduced by Slater [1] in approximation of Hartree-Fock equations. In particular, (1.1) without the Hartree-type is a canonical nonlinear Schrödinger equation, which may model the Bose-Einstein condensate, see e.g., [2]. (1.1) without the power-type is a standard nonlinear Hartree equation, which can be considered as a classical limit of a field equation describing a quantum mechanical nonrelativeistic many-boson system interacting through a two body potential  $|x|^{-2}$ , see e.g., [3]. (1.1) can be viewed as a generalization of the Poission equation for the gravitational potential that is typical in the study of fractional Newtonian gravity as an alterative to standard Newtonian gravity, see e.g.,[4], which has been extensively studied over the last few years, such as [5-7] and the references therein.

For (1.1), the local well-posedness in the natural energy space  $H^1(\mathbb{R}^3)$  is established in [8] (also see [9]). That is, for any initial data  $\psi_0(x) \in H^1(\mathbb{R}^3)$ , there exists a unique maximal solution

$$\psi(t) \in \mathcal{C}((-T_{\min}, T_{\max}), H^1(\mathbb{R}^3)) \cap \mathcal{C}^1((-T_{\min}, T_{\max}), H^{-1}(\mathbb{R}^3))$$

of (1.1) with  $\psi(0) = \psi_0$ . In addition, the mass

$$M(\psi(t)) := ||\psi(t)||_2^2$$

and energy

$$E(\psi(t)) := \frac{1}{2} \|\nabla \psi(t)\|_2^2 - \frac{1}{4} \|(|x|^{-2} * |\psi(t)|^2) |\psi(t)|^2 \|_1 - \frac{1}{p+1} \|\psi(t)\|_{p+1}^{p+1}$$

of the solution are conserved by flow. Moreover, there holds the blow-up criterion

$$T_{\max} < +\infty$$
 implies  $\lim_{t \nearrow T_{\max}} \|\psi(t)\|_{H^1(\mathbb{R}^3)} = +\infty$ .

The sharp threshold of global existence and blowup for (1.1) with general Hartree-type  $(|x|^{-\alpha}*|\psi|^2)\psi$  was considered in [10] for the case  $N\geq 3, 2\leq \alpha<\min\{4,N\}, 1+\frac{4}{N}< p<\frac{N+2}{N-2}$  and in [11,12] for the case  $N\geq 3, 2<\alpha<\min\{4,N\}, 1< p\leq 1+\frac{4}{N}$ . The existence and multiplicity of normalized solutions for (1.1) was obtained in [13]. In particular, the existence and the non-degeneracy of ground state solutions were proved and further the minimal mass blowup solutions were constructed in [4] for (1.1) with  $p=1+\frac{4}{3}$ . However, to our knowledge, the stability and instability (especially strong instability) of standing waves for (1.1) has not been studied in the literature.

By a standing wave, we mean a solution to (1.1) with the form

$$\psi(t,x) = e^{i\omega t} u_{\omega}(x), \qquad \omega \in \mathbb{R}, \ u_{\omega}(x) \in H^1(\mathbb{R}^3) \setminus \{0\},$$