

Existence of Nontrivial Steady State of a Chemotaxis Fluid Coupled System

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Abstract. This paper studied a chemotaxis fluid coupled system which models the so-called “chemotactic Boycott effect”. Under certain inhomogeneous boundary conditions, we proved the existence of nontrivial steady state which depends only on one variable of such system in $[0,1]^3$. A positive lower bound of the oxygen concentration c was also obtained.

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1 Introduction and main results

In this paper, we investigate the steady state of the following coupled chemotaxis fluid model which was firstly proposed in [1]:

$$\begin{cases} n_t + \mathbf{u} \cdot \nabla n - \Delta n = -\nabla \cdot (n\chi(c)\nabla c), & \text{in } \mathbb{R}^+ \times \Omega, \\ c_t + \mathbf{u} \cdot \nabla c - \Delta c = -n\kappa(c), & \text{in } \mathbb{R}^+ \times \Omega, \\ \mathbf{u}_t - \Delta \mathbf{u} + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = -n\nabla \phi, & \text{in } \mathbb{R}^+ \times \Omega, \\ \nabla \cdot \mathbf{u} = 0, & \text{in } \mathbb{R}^+ \times \Omega, \end{cases} \quad (1.1)$$

where c denotes the oxygen concentration, n denotes the bacteria density, and \mathbf{u} denotes the fluid velocity. $\Omega \subset \mathbb{R}^d$ with $d = 2, 3$ is the physical domain. This system models the so-called “chemotactic Boycott effect”, where $\chi(c)$ represents the chemotactic sensitivity and $\nabla \phi$ models the downwards gravitational effect, $\kappa(c)$ models an inactivity threshold

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of the bacteria due to low oxygen supply. Usually(c.f. [2, 3]), $\kappa(c)$ was supposed to be a continuously differentiable, monotonically increasing function with $\kappa(0) = 0$.

The above time-dependent system (1.1) has been extensively studied during the past decade. For the existence of local and global weak solution under certain assumptions on χ and κ , we refer to [3–5]. See [2] for the study of the Cauchy problem of this system. The existence of global large-data solutions was studied in [6]. For the study of stabilization of such system we refer to [7]. For a survey and critical analysis of chemotaxis models, we refer to [8]. More related studies can be found in [9–19] and references therein.

However, to our knowledge, there is no result on nontrivial steady state of (1.1). The only related work we found is [20], in which the authors proved the existence of nonconstant solution to

$$\begin{cases} -\Delta n + \nabla \cdot (n \nabla c) = 0, & \text{in } \Omega, \\ -\Delta c + nc = 0, & \text{in } \Omega, \\ \frac{\partial c}{\partial \nu} = (\gamma - c)g, \quad \frac{\partial n}{\partial \nu} - n \frac{\partial c}{\partial \nu} = 0, & \text{on } \partial \Omega. \end{cases} \quad (1.2)$$

To fill this gap, we consider the following time independent system:

$$\begin{cases} -\Delta n + \mathbf{u} \cdot \nabla n = -\nabla \cdot (n \chi(c) \nabla c), & \text{in } \Omega, \\ -\Delta c + \mathbf{u} \cdot \nabla c = -n \kappa(c), & \text{in } \Omega, \\ -\Delta \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} + \nabla p = -n \nabla \phi, & \text{in } \Omega, \\ \nabla \cdot \mathbf{u} = 0, & \text{in } \Omega, \end{cases} \quad (1.3)$$

where $\Omega = [0, 1]^3$ denotes the unit cubic domain in \mathbb{R}^3 .

1.1 The boundary conditions

In previous studies of chemotaxis fluid system, many authors used the following homogeneous boundary conditions

$$\frac{\partial n}{\partial \nu} = \frac{\partial c}{\partial \nu} = 0, \quad \mathbf{u} = \mathbf{0} \quad \text{on } \partial \Omega. \quad (1.4)$$

It is easy to see that with this boundary condition, the only steady state will be trivial state $(n_0, 0, \mathbf{0})$ with some constant n_0 . In fact, by multiplying both sides of the second equation in (1.3) with c , we get

$$\int_{\Omega} |\nabla c|^2 dx = - \int_{\Omega} n \kappa(c) c dx \leq 0,$$

since n, c are nonnegative and $\kappa(c)$ are always supposed to be increasing with $\kappa(0) = 0$. Thus $c = \text{constant}$. Substituting this fact into the first equation in (1.3), and multiplying both sides of this equation with n , we get

$$\int_{\Omega} |\nabla n|^2 dx = 0.$$